

Concentration of solutions of non-linear elliptic equations involving critical Sobolev exponent

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Abstract

In \mathbb{R}^n ($n \geq 3$), an interesting property of the semi-linear equation

$$\Delta u_o + c_n K_o u_o^{\frac{n+2}{n-2}} = 0$$

is that, when K_o is a positive constant, solutions can concentrate at any point. When K_o is not a constant, we show that concentration of solutions requires strong conditions on K_o . Through the stereographic projection, the discussion can be extended to S^n , and is related to bubbling, or the blow-up phenomenon.

KEY WORDS: Concentration, blow-up, scalar curvature equation.

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§ 1. Introduction

In this article we consider positive smooth solutions u of the equation

$$(1.1) \quad \Delta_{g_1} u - c_n n(n-1)u + c_n K u^{\frac{n+2}{n-2}} = 0 \quad \text{in } S^n.$$

Here Δ_{g_1} is the Laplacian for the standard metric g_1 on the unit sphere S^n ($n \geq 3$), and $c_n = (n-2)/[4(n-1)]$. Equation (1.1) describes the scalar curvature K of the conformal metric $u^{\frac{4}{n-2}} g_1$ [4]. Under the stereographic projection $\mathcal{P} : S^n \rightarrow \mathbb{R}^n$ (cf. [14] and § 7), with

$$(1.2) \quad K_o(y) := K(\mathcal{P}^{-1}(y)), \quad u_o(y) = u(\mathcal{P}^{-1}(y)) \left(\frac{2}{1 + \|y\|^2} \right)^{\frac{n-2}{2}} \quad \text{for } y \in \mathbb{R}^n,$$

equation (1.1) can be expressed as

$$(1.3) \quad \Delta u_o + c_n K_o u_o^{\frac{n+2}{n-2}} = 0 \quad \text{in } \mathbb{R}^n.$$

The geometric accent of the equations is reflected analytically in the critical Sobolev exponent. Together with conformal invariance, they may cause bubbles to appear [25]. Active studies are conducted on existence of solutions and fine asymptotic properties, employing powerful ideas in partial differential equations and global geometry (see, for instances, recent publications [1], [3], [7], [8], [9], [10], [17], [22], and the references within). However, key questions like the Nirenberg problem and the Kazdan-Warner problem remain unresolved.

An exquisite result of Gidas, Ni and Nirenberg ([12], [13]; cf. [5], [23]) shows that when K_o is a positive constant, say (after rescaling), $K_o = 4n(n-1)$, any positive smooth solution of equation (1.3) is of the form

$$(1.4) \quad u_{\lambda,p}(y) := \left(\frac{\lambda}{\lambda^2 + \|y-p\|^2} \right)^{\frac{n-2}{2}} \quad \text{for } y \in \mathbb{R}^n.$$

Here p is a fixed point in \mathbb{R}^n , and λ a positive number. Thus the rigidity and flexibility of the equation are captured. As an interesting consequence, solutions

can concentrate near p when $\lambda \rightarrow 0^+$. Indeed, direct calculation reveals that

$$\begin{aligned} \int_{\mathbb{R}^n} u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy & \text{ is independent on } \lambda \text{ and } p, \\ \int_{\mathbb{R}^n \setminus B_p(\rho)} u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy & \leq C \left(\frac{\lambda}{\rho}\right)^n \longrightarrow 0 \quad \text{as } \lambda \rightarrow 0^+. \end{aligned}$$

In the above, C a positive constant, $B_p(\rho)$ the open ball in \mathbb{R}^n with center at p and radius $\rho > 0$. Observe that $u_{\lambda,p}(p) \rightarrow \infty$ as $\lambda \rightarrow 0^+$. Here we use the term concentration to denote the general phenomenon when the solution is large in a neighborhood of a point, and small outside. The precise meaning is made evident in each theorem.

When K_o is *not* a constant, we show that positive smooth solutions of equation (1.3) in the form of (1.2) can only concentrate on particular places. The first observation is that concentration *cannot* take place at a point p with $K_o(p) \leq 0$ (see propositions 2.2 and 2.4 for the precise statements). Here we make use of the fact that conformal deformations on (S^n, g_1) tend to increase the $L^{\frac{n}{2}}$ -norm of the scalar curvatures [15].

The second obstruction for concentration is $\|\nabla K_o(p)\| \neq 0$, a consequence of the famed Kazdan-Warner *balance* formula:

$$(1.5) \quad \int_{S^n} X(K) u^{\frac{2n}{n-2}} dV_{g_1} = 0.$$

Here X is an arbitrary conformal Killing vector field on (S^n, g_1) (cf. [10] [14]).

Formula (1.5), when projected onto \mathbb{R}^n via \mathcal{P} , and when X is generated by rescaling, gives rise to the Pohozaev identity

$$(1.6) \quad \int_{\mathbb{R}^n} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy = 0.$$

From (1.6), we derive the third obstruction, namely, *high concentration cannot take place at a point p with $\Delta K_o(p) \neq 0$* (the precise statement is found in theorem 4.1). Earlier, Chang-Gursky-Yang [6] and Schoen-Zhang [27] consider similar situation. ((1.6) is satisfied by solutions u_o and K_o related to u and K by (1.2), which guarantees the convergence of the integral and non-existence of boundary terms.

It is not hard to relax this requirement by imposing suitable decay condition on u_o . cf. [6].)

Observe that equation (1.3) is invariant under translations. Using this, we also discover that $\nabla(\Delta K_o)(p) \neq 0$ is an obstruction (theorem 5.1). Further exploration of the Pohozaev identity shows that

$$(1.7) \quad 3 \left(\sum_{i=1}^n \frac{\partial^4 K_o}{\partial y_i^4}(p) \right) + \Delta^2 K_o(p) \neq 0$$

places an additional restriction on strong concentration, see theorem 6.1. Here, $\Delta^2 K_o = \Delta(\Delta K_o)$.

A natural link with the kind of concentrations discussed in this article is found in blow-up or bubbling. Let $\{u_i\} \subset C_+^\infty(S^n)$ be a sequence of solutions of equation (1.1). A point x_b is called a blow-up point of $\{u_i\}$ if there exists a sequence $\{x_i\} \subset S^n$ such that $\lim_{i \rightarrow \infty} u_i(x_i) = \infty$ and $\lim_{i \rightarrow \infty} x_i = x_b$. Point singularity of this type is studied in detail by R. Schoen, Y.Y. Li (cf. [22]), Chen and Lin (cf. [7] [8] [24]), and others. Under suitable conditions [10] [20], u_i can be approximated near x_b by a standard solution as in (1.4).

In order to obtain *a priori* bounds and existence results, methods are developed to eliminate the possibility of blow-up (see the elegant works of Aubin [2], Chang-Gursky-Yang [6], Chen-Li [10], Y.-Y. Li [20] [21], Chen-Lin (op. cit.), Schoen [26], Schoen-Escobar [11], and Schoen-Zhang [27]). Conditions allow, uniform upper bound also implies uniform lower bound, thanks to the Harnack inequality. This becomes crucial as certain blow-up tends to pull down the solutions to zero outside a small neighborhood of the blow-up point x_b (see [8]). The conditions discussed here help to avoid this specific type of concentrations.

Conventions. Throughout this article, $n \geq 3$ is an integer; the functions $u_o \in C_+^\infty(\mathbb{R}^n)$ and $K_o \in C^\infty(\mathbb{R}^n)$ descend from the corresponding functions on S^n via (1.2). We observe the practice of summing over repeated indices, and use C , possibly with sub-indices, to denote various positive constants, which may be rendered differently according to the contents.

§ 2. Zeroth order condition

Let

$$\begin{aligned} V &:= \int_{S^n} u^{\frac{2n}{n-2}} dV_{g_1} = \int_{\mathbb{R}^n} u_o^{\frac{2n}{n-2}} dy, \\ T &:= \int_{S^n} K u^{\frac{2n}{n-2}} dV_1 = \int_{\mathbb{R}^n} K_o u_o^{\frac{2n}{n-2}} dy. \end{aligned}$$

On account of (1.1), we have

$$(2.1) \quad T = c_n^{-1} \int_{S^n} \|\nabla_1 u\|^2 dV_{g_1} + n(n-1) \int_{S^n} |u|^2 dV_{g_1} > 0.$$

Proposition 2.2. *Let K_o be as in (1.2). Assume that $K_o(p) < 0$ for a point $p \in \mathbb{R}^n$. There exist positive constants ρ_o and ε_o such that for any positive smooth solution u_o of equation (1.3) in the form of (1.2), the concentration inequality*

$$(2.3) \quad \int_{\mathbb{R}^n \setminus B_p(\rho)} u_o^{\frac{2n}{n-2}} dy \leq \varepsilon \int_{B_p(\rho)} u_o^{\frac{2n}{n-2}} dy$$

cannot hold for $\rho \leq \rho_o$ and $\varepsilon \leq \varepsilon_o$.

Proof. Take ρ_o to be small enough so that $\sup_{B_p(\rho_o)} K_o < 0$. Set

$$\sigma := - \left(\sup_{B_p(\rho_o)} K_o \right) > 0, \quad \text{and} \quad \varepsilon_o = \left(\sup_{\mathbb{R}^n} K_o \right)^{-1} \sigma.$$

Suppose that (2.3) holds for $\rho \leq \rho_o$ and $\varepsilon \leq \varepsilon_o$. We have

$$\begin{aligned} T &= \int_{B_p(\rho)} K_o u_o^{\frac{2n}{n-2}} dy + \int_{\mathbb{R}^n \setminus B_p(\rho)} K_o u_o^{\frac{2n}{n-2}} dy \\ &\leq -\sigma \int_{B_p(\rho)} u_o^{\frac{2n}{n-2}} dy + \left(\sup_{\mathbb{R}^n} K_o \right) \int_{\mathbb{R}^n \setminus B_p(\rho)} u_o^{\frac{2n}{n-2}} dy \\ &\leq -\sigma \int_{B_p(\rho)} u_o^{\frac{2n}{n-2}} dy + \left(\sup_{\mathbb{R}^n} K_o \right) \varepsilon \int_{B_p(\rho)} u_o^{\frac{2n}{n-2}} dy \leq 0, \end{aligned}$$

which contradicts (2.1). □

From above, it is not immediately clear that concentration cannot take place at a point p with $K_o(p) = 0$. This can be shown with the help of a result in [15].

Proposition 2.4. *Let K_o be as in (1.2). Assume that $K_o(p) = 0$ for a point $p \in \mathbb{R}^n$. Given any positive number C , there exist positive constants ρ_1 and ε_1 such that for any positive smooth solution u_o of equation (1.3) in the form of (1.2), the concentration inequalities*

$$(2.5) \quad \int_{\mathbb{R}^n \setminus B_p(\rho)} u_o^{\frac{2n}{n-2}} dy \leq \varepsilon \int_{B_p(\rho)} u_o^{\frac{2n}{n-2}} dy \quad \text{and} \quad V \leq C$$

do not hold for $\rho \leq \rho_1$ and $\varepsilon \leq \varepsilon_1$.

Proof. By applying lemma 4.5 in [15] on S^n with the standard metric, and then transferring to \mathbb{R}^n by the stereographic projection as in (1.2), we obtain

$$(2.6) \quad \int_{\mathbb{R}^n} |K_o|^{\frac{n}{2}} u_o^{\frac{2n}{n-2}} dy \geq [n(n-1)]^{\frac{n}{2}} \omega_n.$$

Here ω_n is the volume of the standard n -sphere. Take ρ_1 to be small enough so that

$$|K_o(y)|^{\frac{n}{2}} \leq \frac{[n(n-1)]^{\frac{n}{2}} \omega_n}{2C}$$

for $y \in B_p(\rho)$. Let

$$\varepsilon_1 = \frac{[n(n-1)]^{\frac{n}{2}} \omega_n}{2C} \left(\sup_{\mathbb{R}^n} K_o \right)^{-\frac{n}{2}}.$$

Here C is the positive constant in (2.5). Suppose that (2.5) holds for $\rho \leq \rho_1$ and $\varepsilon \leq \varepsilon_1$. We have

$$\begin{aligned} & \int_{\mathbb{R}^n} |K_o|^{\frac{n}{2}} u_o^{\frac{2n}{n-2}} dy \\ &= \int_{B_p(\rho)} |K_o|^{\frac{n}{2}} u_o^{\frac{2n}{n-2}} dy + \int_{\mathbb{R}^n \setminus B_p(\rho)} |K_o|^{\frac{n}{2}} u_o^{\frac{2n}{n-2}} dy \\ &\leq \frac{[n(n-1)]^{\frac{n}{2}} \omega_n}{2C} \int_{B_p(\rho)} u_o^{\frac{2n}{n-2}} dy + \left(\sup_{\mathbb{R}^n} K_o \right)^{\frac{n}{2}} \int_{\mathbb{R}^n \setminus B_p(\rho)} u_o^{\frac{2n}{n-2}} dy \\ &< \frac{[n(n-1)]^{\frac{n}{2}} \omega_n}{2} + \varepsilon V \left(\sup_{\mathbb{R}^n} K_o \right)^{\frac{n}{2}} \\ &= [n(n-1)]^{\frac{n}{2}} \omega_n. \end{aligned}$$

The strict inequality above provides a contradiction with (2.6). \square

Remark 2.7. In proposition 2.4, whether the bound on V can be removed is not known. Interestingly, there are examples which show that V can become very large due to strong concentration at a point, even though K_o is very close to a positive constant (see [18] and [29]; cf. also [16]). However, under mild conditions on K_o (see [17]), it can be shown that if x_b is a blow-up point as defined in the introduction, then $K(x_b) > 0$.

§ 3. First order property

The stereographic projection \mathcal{P} enables us to bring the discussion from \mathbb{R}^n to S^n , or vice versa. For first order obstruction, it is more convenient to consider S^n . Denote by $\mathcal{B}_q(r)$ the open (metric) ball in the standard sphere S^n , where q is the center and $r \in (0, \pi)$ the radius of the ball.

Proposition 3.1. *Let $K \in C^\infty(S^n)$. Assume that $K(q) > 0$ and $\nabla_1 K(q) \neq 0$ for a point $q \in S^n$. There exist positive constants ρ_2 and ε_2 such that for any positive smooth solution u of equation (1.1), the concentration inequality*

$$(3.2) \quad \int_{S^n \setminus \mathcal{B}_q(r)} u^{\frac{2n}{n-2}} dV_{g_1} \leq \varepsilon \int_{\mathcal{B}_q(r)} u^{\frac{2n}{n-2}} dV_{g_1}$$

cannot hold for $r \leq \rho_2$ and $\varepsilon \leq \varepsilon_2$.

Proof. Let $\|\nabla_1 K(q)\| = \delta^2 > 0$. In the Kazdan-Warner formula (1.5), we can choose the coordinate system so that the conformal Killing vector field X has the property that $\|X(q)\| = 1$ and $X(q)$ is in the direction of $\nabla_1 K(q)$. This is possible because of the innate symmetry of S^n . Furthermore, we may take $\|X\| \leq 1$ in S^n . Fix a positive constant ρ_2 such that

$$(3.3) \quad \langle X, \nabla_1 K(x) \rangle_{g_1} \geq \frac{\delta^2}{2} \quad \text{for } x \in \mathcal{B}_q(\rho_2).$$

Let D be a positive constant such that

$$\|\nabla_1 K\| \leq D \quad \text{in } S^n.$$

It follows that $X(K) = \langle X, \nabla_1 K \rangle_{g_1} \leq \|X\| \cdot \|\nabla_1 K\| \leq D$. Take

$$\varepsilon_2 = \frac{\delta^2}{3D}.$$

For $r \leq \rho_2$, we have

$$\begin{aligned} & \int_{S^n} X(K) u^{\frac{2n}{n-2}} dV_{g_1} = 0 \\ \implies & \int_{\mathcal{B}_q(r)} X(K) u^{\frac{2n}{n-2}} dV_{g_1} = - \int_{S^n \setminus \mathcal{B}_q(r)} X(K) u^{\frac{2n}{n-2}} dV_{g_1} \\ \implies & \frac{\delta^2}{2} \int_{\mathcal{B}_q(r)} u^{\frac{2n}{n-2}} dV_{g_1} \leq D \int_{S^n \setminus \mathcal{B}_q(r)} u^{\frac{2n}{n-2}} dV_{g_1} \\ \implies & \int_{S^n \setminus \mathcal{B}_q(r)} u^{\frac{2n}{n-2}} dV_{g_1} \geq \frac{\delta^2}{2D} \int_{\mathcal{B}_q(r)} u^{\frac{2n}{n-2}} dV_{g_1}. \end{aligned}$$

The imbalance renders (3.2) invalid for $\varepsilon \leq \varepsilon_2$ and $r \leq \rho_2$. \square

Related to the above, we refer to [28] and [30] for first order conditions on concentration for certain singularly perturbed elliptic equations in \mathbb{R}^n .

§ 4. Second order property

Because equation (1.3) is invariant under translations, for the moment, we focus the discussion on the origin.

Theorem 4.1. *Let K_o be as in (1.2). Assume that $\nabla K_o(0) = 0$ and $\Delta K_o(0) \neq 0$. Given any positive numbers C and ρ , there exist positive numbers c_1 and c_2 , such that for any positive smooth solution u_o of equation (1.3) in the form of (1.2), the concentration*

$$(4.2) \quad \left\| \left(\frac{u_o}{u_{\lambda,0}} \right)^{\frac{2n}{n-2}} - 1 \right\|_{C^0(B_o(\rho))} \leq \delta, \quad \int_{\mathbb{R}^n \setminus B_o(\rho)} u_o^{\frac{2n}{n-2}} dy \leq C \lambda^3$$

cannot take place for $\lambda \leq c_1$ and $\delta \leq c_2$. Here $u_{\lambda,0}$ is the standard spherical solution defined in (1.4).

Proof. Consider the case $\Delta K_o(0) = \Lambda > 0$ first. As

$$(4.3) \quad \begin{aligned} \int_{\mathbb{R}^n} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy &= 0 \quad (\text{the Pohozaev identity}) \\ \implies \int_{B_o(\rho)} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy &= - \int_{\mathbb{R}^n \setminus B_o(\rho)} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy, \end{aligned}$$

we intend to show that, under concentration as expressed in (4.2), the left hand side of the above is $O(\lambda^2)$, and the other side $O(\lambda^3)$. Thus (4.3) cannot be balanced for small λ . To this end we apply Taylor's expansion and the fact that $\nabla K_o(0) = 0$, obtaining

$$K_o(y) = K_o(0) + \frac{1}{2} \sum_{i,j} y_i y_j \frac{\partial^2 K_o}{\partial y_i \partial y_j} (0) + R(y).$$

Here R is a smooth function with vanishing first and second order derivatives at the origin. It follows that

$$(4.4) \quad y \cdot \nabla K_o(y) = \sum_{i,j} y_i y_j \frac{\partial^2 K_o}{\partial y_i \partial y_j} (0) + y \cdot \nabla R(y).$$

By the remainder theorem for Taylor's expansions, we have

$$(4.5) \quad |y \cdot \nabla R(y)| \leq |y| \cdot \|\nabla R(y)\| \leq C_1 \|y\|^3 \quad \text{for } \|y\| \leq \rho.$$

It follows from (4.4) that

$$(4.6) \quad |y \cdot \nabla K_o(y)| \leq C_2 |y|^2 \quad \text{for } |y| \leq \rho.$$

Assuming that (4.2) holds. We have

$$(4.7) \quad \begin{aligned} &\int_{B_o(\rho)} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy \\ &= \int_{B_o(\rho)} (y \cdot \nabla K_o) u_{\lambda,0}^{\frac{2n}{n-2}} dy + \int_{B_o(\rho)} (y \cdot \nabla K_o) \left[u_o^{\frac{2n}{n-2}} - u_{\lambda,0}^{\frac{2n}{n-2}} \right] dy \\ &\geq \int_{B_o(\rho)} (y \cdot \nabla K_o) u_{\lambda,0}^{\frac{2n}{n-2}} dy - \int_{B_o(\rho)} |y \cdot \nabla K_o| \left| u_o^{\frac{2n}{n-2}} - u_{\lambda,0}^{\frac{2n}{n-2}} \right| dy \\ &\geq \sum_{i,j} \left[\frac{\partial^2 K_o}{\partial y_i \partial y_j} (0) \right] \int_{B_o(\rho)} y_i y_j u_{\lambda,0}^{\frac{2n}{n-2}} dy - C_2 \delta \int_{B_o(\rho)} \|y\|^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy \\ &\quad - C_1 \int_{B_o(\rho)} \|y\|^3 u_{\lambda,0}^{\frac{2n}{n-2}} dy \quad (\text{using (4.2), (4.4) - (4.6)}). \end{aligned}$$

As $u_{\lambda,0}$ depends only on $r = \|y\|$, by symmetry, one obtains

$$\begin{aligned} \int_{B_o(\rho)} y_i y_j u_{\lambda,0}^{\frac{2n}{n-2}} dy &= 0 \quad \text{for } i \neq j, \\ \int_{B_o(\rho)} y_i^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy &= \int_{B_o(\rho)} y_j^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy, \\ \implies \int_{B_o(\rho)} y_i^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy &= \frac{1}{n} \int_{B_o(\rho)} r^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy. \end{aligned}$$

We compute

$$\begin{aligned} (4.8) \quad & \int_{B_o(\rho)} r^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy \\ &= \omega_{n-1} \int_0^\rho \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n r^{n+1} dr \\ &= \omega_{n-1} \int_0^{\arctan(\rho/\lambda)} \left(\frac{\lambda}{\lambda^2 \sec^2 \phi} \right)^n \lambda^{n+2} \tan^{n+1} \phi \sec^2 \phi d\phi \quad (r = \lambda \tan \phi) \\ &= \lambda^2 \omega_{n-1} \int_0^{\arctan(\rho/\lambda)} \cos^{2(n-1)} \phi \left(\frac{\sin^{n+1} \phi}{\cos^{n+1} \phi} \right) d\phi \\ &= \lambda^2 \omega_{n-1} \int_0^{\arctan(\rho/\lambda)} \sin^{n+1} \phi \cos^{n-3} \phi d\phi \quad (n \geq 3) \\ &= \lambda^2 I_{\rho/\lambda}. \end{aligned}$$

Here

$$(4.9) \quad I_{\rho/\lambda} := \omega_{n-1} \int_0^{\arctan(\rho/\lambda)} \sin^{n+1} \phi \cos^{n-3} \phi d\phi,$$

and ω_{n-1} is the volume of the standard sphere S^{n-1} . As it can be seen in (4.9), $I_{\rho/\lambda}$ is bounded from above, and its value is larger for smaller λ , assuming that ρ is fixed. Similarly,

$$(4.10) \quad \int_{B_o(\rho)} \|y\|^3 u_{\lambda,0}^{\frac{2n}{n-2}} dy = \lambda^3 \omega_{n-1} \int_0^{\arctan(\rho/\lambda)} \sin^{n+2} \phi \cos^{n-4} \phi d\phi.$$

When $n \geq 4$,

$$\int_{B_o(\rho)} \|y\|^3 u_{\lambda,0}^{\frac{2n}{n-2}} dy \leq C_4 \lambda^3.$$

When $n = 3$,

$$\begin{aligned}
& \int_0^{\arctan(\rho/\lambda)} \sin^{n+2} \phi \sec \phi d\phi \\
\leq & \int_0^{\arctan(\rho/\lambda)} \sec \phi d\phi \\
= & \ln |\sec y + \tan y|_{y=\arctan(\rho/\lambda)} = \ln \left| \sqrt{1 + \tan^2 y} + \tan y \right|_{y=\arctan(\rho/\lambda)} \\
= & \ln \left(\sqrt{1 + \frac{\rho^2}{\lambda^2}} + \frac{\rho}{\lambda} \right) \leq \ln \left(\frac{3\rho}{\lambda} \right) \leq \sqrt{\frac{3\rho}{\lambda}} \quad (\text{provided } \rho/\lambda \geq 1).
\end{aligned}$$

Thus

$$\int_{B_o(\rho)} \|y\|^3 u_{\lambda,0}^{\frac{2n}{n-2}} dy \leq C_5 \lambda^{\frac{5}{2}} \quad \text{for } \lambda \leq \rho \quad (n \geq 3),$$

where C_5 is a positive constant that depends on ρ and n only. It follows from the symmetry and (4.7)–(4.10) that

$$(4.11) \quad \int_{B_o(\rho)} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy \geq \left[\frac{\Delta K_o(0)}{n} - C \delta \right] I_{\rho/\lambda} \lambda^2 - C_6 \lambda^{\frac{5}{2}}.$$

We choose

$$c_2 = \frac{\Lambda}{2n C}$$

so that when $\delta \leq \Lambda_2$, we have

$$\left[\frac{\Delta K_o(0)}{n} - C \delta \right] \geq \frac{\Lambda}{n} - \frac{\Lambda}{2n} = \frac{\Lambda}{2n}.$$

By the decay property of K_o as in (1.2), there exists a positive constant C_7 such that $|y \cdot \nabla K_o(y)| \leq C_7$ for all $y \in \mathbb{R}^n$. From (4.3) we have

$$\begin{aligned}
(4.12) \quad & \left(I_{\rho/\lambda} \frac{\Lambda}{2n} \right) \lambda^2 - C_6 \lambda^{\frac{5}{2}} \leq \int_{\mathbb{R}^n \setminus B_o(\rho)} |y \cdot \nabla K_o| u_o^{\frac{2n}{n-2}} dy \\
\implies & \left(I_{\rho/\lambda} \frac{\Lambda}{2n} \right) \lambda^2 \leq C_6 \lambda^{\frac{5}{2}} + C_8 \lambda^3 \\
\implies & \frac{\Lambda I_{20}}{2n} \leq C_9 \lambda^{\frac{1}{2}} \quad \left(\text{provided that } \lambda \leq \min \left\{ \frac{\rho}{20}, 1 \right\} \right) \\
\implies & \lambda > \left[\frac{\Lambda I_{20}}{2n C_9} \right]^2.
\end{aligned}$$

Hence we may choose

$$c_1 = \left[\frac{\Lambda I_{20}}{2n C_9} \right]^2.$$

From (4.12), we conclude that (4.2) cannot hold for $\lambda \leq c_1$ and $\delta \leq c_2$. The case $\Delta K_o(0) < 0$ is similar. \square

Remark 4.13. Fixing ρ in (4.2), we observe that when λ is small enough, then (4.2) implies that

$$\int_{\mathbb{R}^n \setminus B_0(\rho)} u_o^{\frac{2n}{n-2}} dy \leq \lambda^2 \int_{B_0(\rho)} u_o^{\frac{2n}{n-2}} dy.$$

(Compare with the calculations in the introduction following (1.4).) It follows that (4.2), when projected back to S^n , also implies inequality of the form (3.2).

§ 5. Third order restriction

Theorem 5.1. *Let K_o be as in (1.2). Assume that $\|\nabla K_o(0)\| = \Delta K_o(0) = 0$ and $\nabla(\Delta K_o)(0) \neq 0$. Given any positive numbers C and ρ , there exist positive constants c_3 and c_4 , such that for any positive smooth solution u_o of equation (1.3) in the form of (1.2), the concentration inequalities*

$$(5.2) \quad \left\| \left(\frac{u_o}{u_{\lambda,0}} \right)^{\frac{2n}{n-2}} - 1 \right\|_{C^0(B_o(\rho))} \leq \delta \lambda, \quad \int_{\mathbb{R}^n \setminus B_o(\rho)} u_o^{\frac{2n}{n-2}} dy \leq C \lambda^3$$

cannot take place for $\lambda \leq c_3$ and $\delta \leq c_4$.

Proof. We proceed as in the proof of theorem 4.1, and observe the effect of translations. Take a point $p = (p_1, \dots, p_n) \in \mathbb{R}^n$ such that

$$(5.3) \quad \gamma := p \cdot \nabla(\Delta K_o)(0) > 0.$$

Consider the translation

$$(5.4) \quad K_p(y) := K_o(y - p) \quad \text{and} \quad u_p := u_o(y - p) \quad \text{for } y \in \mathbb{R}^n.$$

It follows that u_p satisfies the equation

$$(5.5) \quad \Delta u_p + c_n K_p u_p^{\frac{n+2}{n-2}} = 0 \quad \text{in } \mathbb{R}^n.$$

In addition, u_p has similar decay property as expressed in (1.2). We also have

$$(5.6) \quad \left\| \left(\frac{u_p(y)}{u_{\lambda,0}(y-p)} \right)^{\frac{2n}{n-2}} - 1 \right\|_{C^0(B_p(\rho))} \leq \delta \lambda, \quad \int_{\mathbb{R}^n \setminus B_p(\rho)} u_p^{\frac{2n}{n-2}} dy \leq C \lambda^3.$$

Let

$$u_{\lambda,p}(y) := u_{\lambda,0}(y-p) = \left(\frac{\lambda}{\lambda^2 + \|y-p\|^2} \right)^{\frac{n-2}{2}} \quad \text{for } y \in \mathbb{R}^n.$$

One obtains

$$\begin{aligned} & \int_{B_p(\rho)} (y \cdot \nabla K_p) u_p^{\frac{2n}{n-2}} dy \\ &= \int_{B_p(\rho)} [(y-p) \cdot \nabla K_p] u_p^{\frac{2n}{n-2}} dy + \int_{B_p(\rho)} (p \cdot \nabla K_p) u_p^{\frac{2n}{n-2}} dy \\ &\geq \int_{B_p(\rho)} [(y-p) \cdot \nabla K_p] u_{\lambda,p}^{\frac{2n}{n-2}} dy - \int_{B_p(\rho)} |(y-p) \cdot \nabla K_p| \left| u_p^{\frac{2n}{n-2}} - u_{\lambda,p}^{\frac{2n}{n-2}} \right| dy \\ &\quad + \int_{B_p(\rho)} (p \cdot \nabla K_p) u_{\lambda,p}^{\frac{2n}{n-2}} dy - \int_{B_p(\rho)} |p \cdot \nabla K_p| \left| u_p^{\frac{2n}{n-2}} - u_{\lambda,p}^{\frac{2n}{n-2}} \right| dy. \end{aligned}$$

We apply the Taylor expansion

$$\begin{aligned} K_p(y) &= K_p(p) + \frac{1}{2} \sum_{i,j} (y_i - p_i) (y_j - p_j) \frac{\partial^2 K_p}{\partial y_i \partial y_j} (p) \\ &\quad + \frac{1}{3!} \sum_{i,j,k} (y_i - p_i) (y_j - p_j) (y_k - p_k) \frac{\partial^3 K_p}{\partial y_i \partial y_j \partial y_k} (p) + R(y), \\ &= K_o(0) + \frac{1}{2} \sum_{i,j} (y_i - p_i) (y_j - p_j) \frac{\partial^2 K_o}{\partial y_i \partial y_j} (0) \\ &\quad + \frac{1}{3!} \sum_{i,j,k} (y_i - p_i) (y_j - p_j) (y_k - p_k) \frac{\partial^3 K_o}{\partial y_i \partial y_j \partial y_k} (0) + R(y). \end{aligned}$$

Here R is a smooth function with vanishing derivatives (up to at least third order) at p . As in (4.4),

$$(5.7) \quad \begin{aligned} (y-p) \cdot \nabla K_p(y) &= \sum_{i,j} (y_i - p_i) (y_j - p_j) \frac{\partial^2 K_o}{\partial y_i \partial y_j} (0) \\ &\quad + \frac{1}{2} \sum_{i,j,k} (y_i - p_i) (y_j - p_j) (y_k - p_k) \frac{\partial^3 K_o}{\partial y_i \partial y_j \partial y_k} (0) \\ &\quad + (y-p) \cdot \nabla R(y). \end{aligned}$$

By the remainder theorem for Taylor's expansions, we have

$$(5.8) \quad \begin{aligned} |p \cdot \nabla R(y)| &\leq C_1 \|y - p\|^3, \\ |(y - p) \cdot \nabla R(y)| &\leq C_2 \|y - p\|^4 \quad \text{for } \|y - p\| < \rho. \end{aligned}$$

From (5.7), we also have

$$(5.9) \quad |(y - p) \cdot \nabla K_p(y)| \leq C_3 \|y - p\|^2 \quad \text{for } \|y - p\| < \rho.$$

Likewise,

$$(5.10) \quad \begin{aligned} &p \cdot \nabla K_p(y) \\ &= \frac{1}{2} \sum_{i,j} [p_i (y_j - p_j) + p_j (y_i - p_i)] \frac{\partial^2 K_o}{\partial y_i \partial y_j} (0) \\ &\quad + \frac{1}{3!} \sum_{i,j,k} [p_i (y_j - p_j) (y_k - p_k) + (y_i - p_i) p_j (y_k - p_k) \\ &\quad \quad + (y_i - p_i) (y_j - p_j) p_k] \frac{\partial^3 K_o}{\partial y_i \partial y_j \partial y_k} (0) + p \cdot \nabla R(y). \end{aligned}$$

By symmetry,

$$\begin{aligned} &\int_{B_p(\rho)} (y_j - p_j) u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy = 0, \\ &\int_{B_p(\rho)} (y_i - p_i) (y_j - p_j) u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy = 0 \quad \text{for } i \neq j, \\ &\int_{B_p(\rho)} (y_i - p_i)^2 u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy = \frac{1}{n} \int_{B_p(\rho)} r^2 u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy \quad (\text{here } r^2 = \|y - p\|^2), \\ &\int_{B_p(\rho)} (y_i - p_i) (y_j - p_j) (y_k - p_k) u_{\lambda,p}^{\frac{2n}{n-2}}(y) dy = 0 \quad \text{for } 1 \leq i, j, k \leq n. \end{aligned}$$

Similar to the proof of theorem 4.1, $\Delta K_p(p) = \Delta K_o(0) = 0$ and (5.8) imply that

$$\begin{aligned} \int_{B_p(\rho)} [(y - p) \cdot \nabla K_p] u_{\lambda,p}^{\frac{2n}{n-2}} dy &= \int_{B_p(\rho)} [(y - p) \cdot \nabla R] u_{\lambda,p}^{\frac{2n}{n-2}} dy \\ &\leq C \lambda^4 \omega_{n-1} \int_0^{\arctan(\rho/\lambda)} \sin^{n+3} \phi \cos^{n-5} \phi d\phi. \end{aligned}$$

In the above, if $n = 3$, we have

$$\int_0^{\arctan(\rho/\lambda)} \sin^{n+3} \phi \cos^{n-5} \phi d\phi \leq \int_0^{\arctan(\rho/\lambda)} \sec^2 \phi d\phi = \rho/\lambda.$$

$$\int_{B_p(\rho)} |p \cdot \nabla K_p| \left| u_p^{\frac{2n}{n-2}} - u_{\lambda,p}^{\frac{2n}{n-2}} \right| dy \leq C_6 \delta \lambda^2,$$

where we use (5.6) and the estimate $|p \cdot \nabla K_p| \leq C_7 \|y - p\|$ for $y \in B_p(\rho)$. Using (4.3), (5.10), (5.12) and (5.13) (compare also with the proof of theorem 4.1), we obtain a contradiction when δ and λ are small enough. \square

§ 6. Fourth order

One may ask what is likely to happen when $\Delta K(0) = \|\nabla K(0)\| = 0$? (Interesting examples include *homogeneous harmonic polynomials* of higher degrees, see the next section.) The method expounded in theorem 4.1 can be used to search for algebraic relations on higher order derivatives of K , cf. [20]. Here we continue with the fourth order condition.

Theorem 6.1. *For $n \geq 5$, Let K_o be as in (1.2). Assume that $\|\nabla K_o(0)\| = \Delta K(0) = 0$, and*

$$(6.2) \quad \Upsilon := 3 \left(\sum_{i=1}^n \frac{\partial^4 K_o}{\partial y_i^4}(0) \right) + \Delta^2 K_o(0) \neq 0.$$

Given positive constants C and ρ , there exist positive numbers c_5 and c_6 such that for any positive smooth solution u_o of equation (1.3) in the form of (1.2), the concentration inequalities

$$(6.3) \quad \left\| \left(\frac{u_o}{u_{\lambda,0}} \right)^{\frac{2n}{n-2}} - 1 \right\|_{B_o(\rho)} \leq \delta \lambda^2, \quad \int_{\mathbb{R}^n \setminus B_o(\rho)} u_o^{\frac{2n}{n-2}} dy \leq C \lambda^5$$

cannot hold for $\lambda \leq c_5$ and $\delta \leq c_6$.

Proof. We explore the main ideas in the proof of theorem 4. Starting with the case that $\Upsilon > 0$, consider the following Taylor expansion

$$\begin{aligned} K_o(y) &= K_o(0) + \frac{1}{2} \sum_{i,j} y_i y_j \frac{\partial^2 K_o}{\partial y_i \partial y_j}(0) + \frac{1}{3!} \sum_{i,j,k} y_i y_j y_k \frac{\partial^3 K_o}{\partial y_i \partial y_j \partial y_k}(0) \\ &\quad + \frac{1}{4!} \sum_{i,j,k,l} y_i y_j y_k y_l \frac{\partial^4 K_o}{\partial y_i \partial y_j \partial y_k \partial y_l}(0) + R_5(y), \end{aligned}$$

which implies that

$$\begin{aligned} y \cdot \nabla K_o(y) &= \sum_{i,j} y_i y_j \frac{\partial^2 K_o}{\partial y_i \partial y_j}(0) + \frac{1}{2!} \sum_{i,j,k} y_i y_j y_k \frac{\partial^3 K_o}{\partial y_i \partial y_j \partial y_k}(0) \\ &\quad + \frac{1}{3!} \sum_{i,j,k,l} y_i y_j y_k y_l \frac{\partial^4 K_o}{\partial y_i \partial y_j \partial y_k \partial y_l}(0) + y \cdot \nabla R_5(y). \end{aligned}$$

Here

$$(6.4) \quad |y \cdot \nabla R_5(y)| \leq C_1 |y|^5 \quad \text{for } |y| < \rho.$$

As in the proof of theorem 4.1, by symmetry and the fact that $\Delta K_o(0) = 0$, we have

$$\begin{aligned} \sum_{i,j} \left(\frac{\partial^2 K_o}{\partial y_i \partial y_j}(0) \int_{B_o(\rho)} y_i y_j u_{\lambda,0}^{\frac{2n}{n-2}} dy \right) &= 0, \\ \int_{B_o(\rho)} y_i y_j y_k u_{\lambda,0}^{\frac{2n}{n-2}} dy &= 0 \quad \text{for } 1 \leq i, j, k \leq n, \\ \int_{B_o(\rho)} y_i y_j y_k y_l u_{\lambda,0}^{\frac{2n}{n-2}} dy &= 0 \quad \text{for } i, j, k, l \text{ being all distinct,} \\ \int_{B_o(\rho)} y_i^2 y_j y_k u_{\lambda,0}^{\frac{2n}{n-2}} dy &= 0 \quad \text{for } j \neq k, \\ \int_{B_o(\rho)} y_i^3 y_j u_{\lambda,0}^{\frac{2n}{n-2}} dy &= 0 \quad \text{for } i \neq j. \end{aligned}$$

Assuming that (6.3) holds, it follows as in (4.7) that

$$\begin{aligned} (6.5) \quad &\int_{B_o(\rho)} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy \\ &\geq \frac{1}{3!} \left[\sum_i \frac{\partial^4 K_o}{\partial y_i^4}(0) \int_{B_o(\rho)} y_i^4 u_{\lambda,0}^{\frac{2n}{n-2}} dy + \sum_{i \neq j} \frac{\partial^4 K_o}{\partial y_i^2 \partial y_j^2}(0) \int_{B_o(\rho)} y_i^2 y_j^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy \right] \\ &\quad - C_2 \delta \int_{B_o(\rho)} \lambda^2 |y|^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy - C_3 \int_{B_o(\rho)} |y|^5 u_{\lambda,0}^{\frac{2n}{n-2}} dy. \end{aligned}$$

To compute the first two integrals in (6.5), let θ be the angle to the y_n -axis. That is, $y_n = |y| \cos \theta$. Set

$$\begin{aligned} I_n &:= \omega_{n-2} \int_0^\pi \cos^4 \theta \sin^{n-2} \theta d\theta = \frac{3 \omega_{n-1}}{n(n+2)}, \\ J_n &:= \frac{\omega_{n-2}}{n-1} \int_0^\pi \cos^2 \theta \sin^n \theta d\theta = \frac{\omega_{n-1}}{n(n+2)} \quad \implies \quad I_n = 3 J_n. \end{aligned}$$

Here we use the formulas

$$\begin{aligned}\int_0^\pi \sin^l \theta \cos^m \theta d\theta &= \frac{m-1}{m+l} \int_0^\pi \sin^l \theta \cos^{m-2} \theta d\theta, \\ \omega_{n-1} &= \omega_{n-2} \int_0^\pi \sin^{n-2} \theta d\theta,\end{aligned}$$

where $l \geq 1$, $m \geq 2$, $n \geq 3$, and ω_{n-2} is the volume of the standard sphere S^{n-2} .

It follows that

$$\begin{aligned}\int_{B_o(\rho)} y_n^4 u_{\lambda,0}^{\frac{2n}{n-2}} dy &= \int_0^\rho \left[\omega_{n-2} \int_0^\pi (r \cos \theta)^4 (r \sin \theta)^{n-2} r d\theta \right] \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n dr \\ &= \omega_{n-2} \int_0^\rho \left[\int_0^\pi \cos^4 \theta \sin^{n-2} \theta d\theta \right] \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n r^{n+3} dr \\ &= I_n \int_0^\rho \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n r^{n+3} dr \\ &= I_n \lambda^4 \int_0^{\arctan(\rho/\lambda)} \sin^{n+3} \phi \cos^{n-5} \phi d\phi \\ &\quad \text{(using the substitution } r = \lambda \tan \phi \text{)}.\end{aligned}$$

On the other hand,

$$\begin{aligned}&\int_{B_o(\rho)} y_n^2 (y_1^2 + \cdots + y_{n-1}^2) u_{\lambda,0}^{\frac{2n}{n-2}} dy \\ &= \int_{B_o(\rho)} y_n^2 (r^2 - y_n^2) u_{\lambda,0}^{\frac{2n}{n-2}} dy \\ &= \int_0^\rho \left[\omega_{n-2} \int_0^\pi (r \cos \theta)^2 (r \sin \theta)^2 (r \sin \theta)^{n-2} r d\theta \right] \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n dr \\ \implies &(n-1) \int_{B_o(\rho)} y_i^2 y_j^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy \quad (i \neq j) \\ &= \omega_{n-2} \int_0^\rho \left(\int_0^\pi \cos^2 \theta \sin^2 \theta \sin^{n-2} \theta d\theta \right) \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n r^{n+3} dr \\ \implies &\int_{B_o(\rho)} y_i^2 y_j^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy = J_n \int_0^\rho \left(\frac{\lambda}{\lambda^2 + r^2} \right)^n r^{n+3} dr \\ \implies &\int_{B_o(\rho)} y_i^2 y_j^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy = J_n \lambda^4 \int_0^{\arctan(\rho/\lambda)} \sin^{n+3} \phi \cos^{n-5} \phi d\phi, \quad i \neq j.\end{aligned}$$

Setting $\sigma = \Upsilon J_n$, we obtain

$$\begin{aligned}
(6.7) \quad & \int_{B_o(\rho)} (y \cdot \nabla K_o) u_o^{\frac{2n}{n-2}} dy \\
& \geq \left(\frac{\sigma}{6}\right) \lambda^4 \int_0^{\arctan(\rho/\lambda)} \sin^{n+3} \phi \cos^{n-5} \phi d\phi - C_2 \delta \int_{B_o(\rho)} \lambda^2 \|y\|^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy \\
& \quad - C_3 \int_{B_o(\rho)} \|y\|^5 u_{\lambda,0}^{\frac{2n}{n-2}} dy.
\end{aligned}$$

As in (4.8) and (4.9),

$$\begin{aligned}
(6.8) \quad & \int_{B_o(\rho)} \lambda^2 \|y\|^2 u_{\lambda,0}^{\frac{2n}{n-2}} dy = \lambda^4 \int_0^{\arctan(\rho/\lambda)} \sin^{n+1} \phi \cos^{n-3} \phi \leq C_4 \lambda^4, \\
& \int_{B_o(\rho)} \|y\|^5 u_{\lambda,0}^{\frac{2n}{n-2}} dy = \lambda^5 \int_0^{\arctan(\rho/\lambda)} \sin^{n+4} \phi \cos^{n-6} \phi \leq C_5 \lambda^{5-\varepsilon}.
\end{aligned}$$

Here $n \geq 5$ and $\varepsilon \in (0, 1)$ is a positive constant. With (6.3), (6.7) and (6.8), we conclude as in the proof of theorem 4.1 that contradiction arises when λ and δ are small enough. The case $\Upsilon < 0$ is similar. \square

§ 7. Homogeneous harmonic polynomials

Here we present some simple functions which satisfy the conditions in theorem 6.1. Let

$$(7.1) \quad Q_k(x) = \sum C_{i_1, \dots, i_k} x_{i_1} \cdots x_{i_k}$$

be a *homogeneous harmonic polynomial* of degree $k \geq 2$ in \mathbb{R}^{n+1} . It is shown in [19] that Q_k satisfies the Kazdan-Warner type identity, namely,

$$\int_{S^n} X(Q_k) dV_{g_1} = 0$$

for any conformal Killing vector field X on S^n .

Assuming that the indices i_1, \dots, i_k in (7.1) are all smaller than $n+1$. Consider the stereographic projection \mathcal{P} from

$$S^n := \{x = (x_1, \dots, x_{n+1}) \in \mathbb{R}^{n+1} \mid \|x\|^2 = 1\}$$

to \mathbb{R}^n , with Cartesian coordinates (y_1, \dots, y_n) . It is given by

$$(7.2) \quad \begin{aligned} y_i &= \frac{x_i}{1 - x_{n+1}}, \quad 1 \leq i \leq n, \\ x_i &= \frac{2y_i}{1 + \|y\|^2}, \quad 1 \leq i \leq n, \quad \text{and} \quad x_{n+1} = \frac{\|y\|^2 - 1}{\|y\|^2 + 1}. \end{aligned}$$

Using \mathcal{P} , we transfer Q_k into \mathbb{R}^n and obtain

$$(7.3) \quad Q_k(y) = \left(\frac{2}{1 + \|y\|^2} \right)^k \sum C_{i_1, \dots, i_k} y_{i_1} \cdots y_{i_k}.$$

It follows that $\nabla Q_k(0) = 0$ (as $k \geq 2$). Moreover,

$$\begin{aligned} \Delta_y Q_k(y) &= \left(\frac{2}{1 + \|y\|^2} \right)^k \Delta_y \left(\sum C_{i_1, \dots, i_k} y_{i_1} \cdots y_{i_k} \right) \\ &\quad - \frac{2^{k+1} k}{(1 + \|y\|^2)^{k+1}} \sum C_{i_1, \dots, i_k} [y \cdot \nabla_y (y_{i_1} \cdots y_{i_k})] \\ &\quad + \left[\frac{2^{k+2} k(k+1) \|y\|^2}{(1 + \|y\|^2)^{k+2}} - \frac{2^{k+1} k n}{(1 + \|y\|^2)^{k+1}} \right] \sum C_{i_1, \dots, i_k} y_{i_1} \cdots y_{i_k} \\ &= 0 - \frac{2^{k+1} k^2}{(1 + \|y\|^2)^{k+1}} \sum C_{i_1, \dots, i_k} y_{i_1} \cdots y_{i_k} \\ &\quad + \left[\frac{2^{k+2} k(k+1) \|y\|^2}{(1 + \|y\|^2)^{k+2}} - \frac{2^{k+1} k n}{(1 + \|y\|^2)^{k+1}} \right] \sum C_{i_1, \dots, i_k} y_{i_1} \cdots y_{i_k} \\ &= 2^{k+1} k \left[\frac{(k+2-n) \|y\|^2 - (n+k)}{(1 + \|y\|^2)^{k+2}} \right] Q_k(y) \\ \implies \quad \Delta_y Q_k(0) &= 0. \end{aligned}$$

Likewise, $\nabla_y (\Delta_y Q_k)(0) = 0$ and $\Delta_y (\Delta_y Q_k)(0) = 0$. Here we make use of the fact that Q_k is a harmonic polynomial and x_{n+1} is not present in $Q_k(x)$.

Using above, one can construct the desired functions. For instance, let Q_4 be the homogeneous harmonic polynomial defined by

$$Q_4(x) := x_1^4 + x_2^4 - 6x_1^2 x_2^2 \quad \text{for } x \in S^n \subset \mathbb{R}^{n+1}.$$

Using the stereographic projection, we obtain

$$Q_4(y) = \left(\frac{2}{1 + \|y\|^2} \right)^4 (y_1^4 + y_2^4 - 6y_1^2 y_2^2).$$

Let $K_o := 1 + Q_4$ in \mathbb{R}^n . We have

$$K_o(0) = 1 > 0, \quad \nabla K_o(0) = 0, \quad \Delta K_o(0) = 0, \quad \nabla(\Delta K_o)(0) = 0,$$

but
$$3 \left(\sum_{i=1}^n \frac{\partial^4 Q_4}{\partial y_i^4}(0) \right) + \Delta_y^2 Q_4(0) = 3 \left(\sum_{i=1}^n \frac{\partial^4 Q_4}{\partial y_i^4}(0) \right) \neq 0.$$

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